# Variational Approach to Protein Design and Extraction of Interaction Potentials 

Flavio Seno, ${ }^{1}$ Cristian Micheletti, ${ }^{2}$ Amos Maritan, ${ }^{2}$ and Jayanth R. Banavar ${ }^{3}$<br>${ }^{1}$ Dipartimento di Fisica and INFM, Università di Padova, Via Marzolo 8, 35131 Padova, Italy<br>${ }^{2}$ International School for Advanced Studies, (S.I.S.S.A. and INFM), Via Beirut 2-4, 34014 Trieste, Italy and The Abdus Salam Centre for Theoretical Physics, Trieste, Italy<br>${ }^{3}$ Department of Physics and Center for Materials Physics, The Pennsylvania State University, 104 Davey Laboratory, University Park, Pennsylvania 16802

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#### Abstract

We present and discuss a novel approach to the direct and inverse protein folding problem. The proposed strategy is based on a variational approach that allows the simultaneous extraction of amino acid interactions and the low-temperature free energy of sequences of amino acids. The knowledgebased technique is simple and straightforward to implement even for realistic off-lattice proteins because it does not entail threadinglike procedures. Its validity is assessed in the context of a lattice model by means of a variety of stringent checks. [S0031-9007(98)07051-3]


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Two long-standing challenges in molecular biology are the direct and inverse problems of protein folding [1]. The first deals with the determination of the thermodynamically stable conformation of a known sequence of amino acids $[2,3]$ while the second involves the elucidation of the amino acid sequence (if any) which admits a given target structure as its stable conformation [4-9]. One route to a solution of the direct problem requires knowledge of the interaction potentials between the protein constituents-in principle one studies the energies of the sequence in various conformations and identifies the native state structure as being the one with the lowest energy [2]. A solution of the inverse or the sequence design problem entails knowledge of the free energies of the sequences [5]. This follows from an application of Boltzmann statistics - what matters is not how low the energy of a sequence is in the target conformation (a measure of which can be obtained from the knowledge of the interaction potentials) but whether this energy is lower than those in alternative conformations.

In this Letter, we introduce a variational approach for extracting the interaction potentials between the protein constituents and the free energies of candidate sequences simultaneously. The method is general and applicable to real proteins. The input is a set of sequences and their respective native structures, as available from the protein data bank (PDB). A feature of the technique is that alternative conformations that compete with the native state in housing a given sequence are not required. Here we apply this method to a lattice model of proteins to provide a stringent validation of the approach. Unlike real proteins, the interaction potentials in such a model system can be chosen and one may measure the accuracy with which these potentials are determined as well as the effectiveness of sequence design.

We adopt the approach of treating proteins at a mesoscopic level with the amino acids being the fundamental units. The influence of the internal degrees of freedom
associated with each amino acid as well as the solvent degrees of freedom is incorporated through a coarse-grained Hamiltonian with effective interactions between the amino acids.

The free energy, $F(S)$, of a sequence, $S$, is defined from the equation

$$
\begin{equation*}
e^{-\beta F(S)}=\sum_{\Gamma} e^{-\beta \mathcal{H}(S, \Gamma)} \tag{1}
\end{equation*}
$$

where $\mathcal{H}(S, \Gamma)$ is the energy of $S$ mounted on a conformation $\Gamma$ and the sum is taken over all conformations that the sequence can adopt. A rigorous solution [5] of the design problem on a target structure $\Gamma$ entails the identification of the sequence(s), $S$, that maximizes the functional

$$
\begin{equation*}
P_{\Gamma}(S)=e^{-\beta[\mathcal{H}(S, \Gamma)-F(S)]} \tag{2}
\end{equation*}
$$

evaluated at a low temperature [below the folding transition temperature where $\left.P_{\Gamma}(S)=1 / 2\right] . \quad P_{\Gamma}(S)$ is the probability that a sequence $S$ is found in conformation $\Gamma$ at an inverse temperature $\beta$. Thus the solution $S$ is the sequence which has the highest low-temperature probability of being found on $\Gamma$. At low temperatures [10], a sequence $\bar{S}$ with a unique ground state, $\bar{\Gamma}$, satisfies the inequality

$$
\begin{equation*}
H(\bar{S}, \bar{\Gamma})-F(\bar{S}) \leq H(S, \bar{\Gamma})-F(S) \tag{3}
\end{equation*}
$$

for arbitrary sequence, $S$, with the equality possibly holding only when $S$ admits $\bar{\Gamma}$ as its native state. A range of such equalities could be used to determine optimal values of variational parameters characterizing the interactions and the low-temperature free energies.

The maximization of $P_{\Gamma}(S)$ is computationally demanding because it involves the calculation of $F(S)$ for each amino acid sequence and an exact calculation of $F(S)$ for a given sequence $S$ involves a sum over the enormous number of its possible conformations. The use of importance-sampling techniques for the estimation of $F(S)$ at low $T$ requires efficient algorithms to find conformations that compete significantly with $\Gamma$ [5]. Such an
approach has been used fruitfully for lattice models of proteins [8] but is not feasible for realistic off-lattice cases [9].
$F(S)$ formally depends only on $S$ and hence one may postulate a functional form of $F$ which depends on sequence properties (e.g., the concentration of amino acids) $[8,9]$. At $T=0$ [10], the free energy of a sequence ought to be exactly equal to its energy in the native state conformation (which depends on the conformation and the interaction potentials) - this forms the basis for our variational approach. Unlike the inequalities (3), the new approach does not entail the mounting of a sequence on any but its own native state conformation. We define an intensive functional $\Delta$ (whose choice is not unique), whose minimization can be used to identify a consistent set of potential and free energy parameters. A convenient choice that we used in our calculations is

$$
\begin{align*}
\Delta= & \bar{\epsilon}\left\{\sum_{i}\left(\frac{\mathcal{H}\left(S_{i}, \Gamma_{i}\right)-F\left(S_{i}\right)}{L_{i}}\right)^{2}\right. \\
& \left.+\left(\frac{\mathcal{H}\left(S_{i}, \Gamma_{i}\right)-F\left(S_{i}\right)}{L_{i}}\right)^{4} \Theta\left[F\left(S_{i}\right)-\mathcal{H}\left(S_{i}, \Gamma_{i}\right)\right]\right\} \tag{4}
\end{align*}
$$

where $\Theta[x]$ is the Heaviside function, and the sum is taken over the sequence-native state conformation set in the protein data bank, and $L_{i}$ is the length of the $i$ th sequence. The second term in (4) is used to penalize cases for which the parameters violate the physical constraint, $\mathcal{H}\left(S_{i}, \Gamma_{i}\right) \geq F\left(S_{i}\right)$. The quantity $\bar{\epsilon}$ is the absolute value of the average of the interaction strengths between amino acids and its utility is explained below. A zero value for $\Delta$ would correspond to a perfect parametrization of both the interaction potentials and the free energies for the finite set of sequences in the data bank. More generally, for a finite protein data bank, there will exist a nonzero region in the parameter space of potentials and free energies within which $\Delta$ is at a minimum. With perfect parametrization, this region would be expected to shrink around the parameter values as the data bank size increases [11]. It should be stressed that, contrary
to common potential extraction or design procedures, the minimization of the functional (4) does not involve the use of decoy structures nor the mounting of sequence $i$ on any structure other than its ground state, $\Gamma_{i}$. In order to create a data bank, a random exploration of the ensemble of 4 -amino-acid sequences of length 16 was performed to select those admitting a unique ground state conformation. The possible protein conformations were assumed to be self-avoiding oriented walks embedded on a square lattice [12] with an interaction between amino acids $i$ and $j$ only if they are next to each other on the lattice and yet not next to each other along the sequence.

We chose an interaction matrix, $\boldsymbol{\epsilon}$, between the four different types (or classes) of amino acids. These are the entries of the $4 \times 4 \epsilon$ matrix in the first row of Table I (with $\epsilon_{1,1}=-40$ ).
To mimic the thermodynamic stability of proteins, we further selected the sequences and retained only those with an energy gap between the unique native state and the first excited state energies $\geq 10$, a constraint satisfied by, roughly, $1 \%$ of the sequences. Our final data bank consisted of 500 sequences with their ground state, $\left\{S_{i}, \Gamma_{i}\right\}_{i=1, \ldots, 500}$.

In our model studies, we chose to parametrize the interaction matrix with the same functional form as the true interaction matrix but with nine variational parameters in the symmetric $\epsilon$ matrix ( $\epsilon_{1,1}$ was held fixed at a value of -40 in order to set the energy scale). We assumed the simplest form for an extensive free energy [8,9] with four variational parameters (denoted by $a_{i}$, $i=1 \ldots 4)$ :

$$
\begin{equation*}
F(S)=a_{1} n_{1}+a_{2} n_{2}+a_{3} n_{3}+a_{4} n_{4}, \tag{5}
\end{equation*}
$$

where $n_{i}$ is the number of amino acids of type $i$ found in $S$. Equation (5) may be viewed as the lowest order expansion of $F$ in the "order parameters," $n_{i}$ 's.
$\Delta$ was minimized using a simulated annealing procedure by constraining the interaction energies, $\epsilon_{i, j}$, to satisfy the hierarchy of strengths deduced from the

TABLE I. A summary of the results with two data banks containing 500 and 250 training proteins, respectively. In all cases $\epsilon_{1,1}$ was fixed at -40 in order to set the energy scale. The row entitled TRUE shows the true potential parameters in both cases. The other rows show the values of the extracted parameters of the potential and the free energy with the number in the first column showing the number of proteins in the training set. A single randomly chosen set was employed in each case. For the second data bank, the folding success rate was $91 \%$, while the unique and degenerate design success rates were $73 \%$ and $96 \%$, respectively.

| PDB Size | $\epsilon_{1,2}$ | $\epsilon_{1,3}$ | $\epsilon_{1,4}$ | $\epsilon_{2,2}$ | $\epsilon_{2,3}$ | $\epsilon_{2,4}$ | $\epsilon_{3,3}$ | $\epsilon_{3,4}$ | $\epsilon_{4,4}$ | $a_{1}$ | $a_{2}$ | $a_{3}$ | $a_{4}$ |
| :---: | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| TRUE | -30 | -20 | -17 | -25 | -13 | -10 | -5 | -2 | -1 |  |  |  |  |
| 100 | -32.63 | -26.80 | -22.71 | -31.57 | -22.71 | -17.76 | -17.76 | -17.75 | -0.00 | -26.50 | -17.44 | -12.60 | -10.39 |
| 200 | -32.62 | -25.56 | -23.17 | -30.65 | -23.17 | -16.06 | -14.07 | -5.43 | -0.00 | -26.15 | -17.14 | -12.33 | -11.07 |
| 300 | -31.27 | -27.91 | -23.29 | -27.91 | -23.17 | -12.61 | -8.50 | -8.50 | -0.00 | -27.17 | -15.62 | -11.83 | -10.87 |
| 400 | -31.03 | -27.14 | -23.27 | -27.14 | -22.44 | -12.50 | -6.94 | -6.77 | -0.00 | -27.62 | -15.25 | -11.52 | -10.38 |
| 500 | -32.23 | -25.93 | -24.12 | -28.55 | -22.78 | -16.40 | -11.79 | -9.13 | -5.21 | -25.63 | -16.49 | -10.53 | -9.55 |
| TRUE | -22 | -18 | -12 | -11 | -17 | -1 | -28 | -13 | -1 |  |  |  |  |
| 250 | -24.05 | -17.95 | -13.22 | -13.02 | -17.95 | -0.00 | -24.06 | -13.22 | -0.00 | -26.26 | -10.41 | -12.70 | -4.14 |

frequencies of pair contacts in the data bank (the more frequent, the stronger). This allowed for a restriction of the search in parameter space.

The quantity $\overline{\boldsymbol{\epsilon}}$ in (4) was useful in avoiding convergence to a spurious trivial solution in which all the $\epsilon_{i, j}$ 's are equal to $\epsilon_{1,1}=-40$, and $F(S)$ becomes $(-40)$ times the number of contacts.

The functional $\Delta$ was minimized using subsets of our global data bank within which the number of elements ranged from 100 to 500 . The minimization was carried out using a simulated annealing algorithm. On the average, for each elementary move, $1 / 4$ of the parameters in $\mathcal{H}$ and $F$ were varied simultaneously by adding to each an independent random quantity picked in the interval $[-\delta,+\delta]$. At the beginning $\delta$ was taken to be of order unity and was then decreased proportionally with the annealing temperature. The temperature was reduced in the annealing process by steps of $5 \%$ with the system being equilibrated at each temperature.

The extracted potentials, as well as the free energy coefficients appear in Table I. We further checked, using the extracted parameters, whether each sequence in the data bank recognized the associated structure as its ground state among all the possible conformations. The success rate was typically $>80 \%$ with an increase in the success rate on increasing the size of the data bank.

We then proceeded to use the functional $(H-F)$ to carry out the sequence design on a target structure. This entails the identification from among the $4^{16}$ sequences the one that minimizes $(H-F)$ (using the extracted parameters) on the target structure $\Gamma$. The correctness of the design is checked by using the true Hamiltonian to verify whether the designed sequence admits $\Gamma$ as its (possibly degenerate) ground state. Our test was performed on 100 structures taken from our data bank using a Monte Carlo procedure.

Figure 1 shows a plot of the design success rate as a function of the size of the training set. It is worth noting that none of the designed sequences appeared in the original data bank. Our analysis was not limited to those sequences with the lowest $(H-F)$ score; we extended it to the ten highest ranking sequences for each target structure. Using the parameters deduced from the training set of size 500 , we found an excellent overall design success of $88 \%$ and $92 \%$ for unique and degenerate encoding, respectively.

In Fig. 2 we have plotted the histogram of the $(H-F)$ distribution for the improperly chosen sequences (black) and the correct ones (gray). The $(H-F)$ score for the improperly chosen sequences takes on large positive values, signaling that the estimated energy $F$ of the sequence in its unknown native state is substantially lower than in the target structure. Thus one may discard a priori the majority of bad solutions by a mere inspection of their large $(H-F)$ scores. The unphysical negative values of $(H-F)$ originate from the small size of the training set and the imperfect parametrization of the free energy. We


FIG. 1. Plot of the success rate in identifying the sequence that admits a preassigned target structure as its degenerate (squares) and nondegenerate (circles) ground state as a function of the training set size. The results were obtained with a single randomly chosen set of each size.
also considered several generalizations of (5) including two-body terms of the form $n_{i} n_{j}$ and chemical potentials that control the number of "walls" separating segments of identical amino acids [9] with slight improvement in the success rates. A further check of the quality of the designed sequences was performed by inspecting the distribution of their energy gaps versus those used in the data bank. The designed sequences tend to have energy gaps between the native state and the first excited state that are larger than those of sequences in the data bank (Fig. 3) showing that the design procedure yields sequences with a higher thermodynamic stability.

Finally, we performed a challenging blind test to assess the validity of the variational approach. The coefficients extracted for the 16 -bead case were used to carry out a sequence design on a compact target conformation of length 25 . The target conformation was $R R D L D R D D L U L D L L U R U L U U R D R D$ (where $R, L$, $U$, and $D$ stand for right, left, up, and down, respectively,


FIG. 2. Distribution of the quantity $(H-F)$ for the correctly chosen sequences (gray) and the improper sequences (black).


FIG. 3. Distribution of the energy gaps between the native state and the first excited state energies for the sequences in the data bank (gray) and designed sequences that have a nondegenerate ground state (black).
and indicate the directions of the bonds that define the self-avoiding lattice conformation) which is highly designable [13] in a different lattice model [14] and is geometrically regular [13]. The sequence chosen was 124211324211324211324211 . Indeed, an exhaustive search of the native state of the sequence among all selfavoiding walks of length 25 confirmed that this sequence had the target structure as its unique ground state.

In order to ensure that the strategy used here is robust and independent of the particular choice of the $\epsilon$ matrix and/or data bank, we performed a similar analysis using another randomly generated interaction matrix and found results of statistically similar quality as summarized in the bottom of Table I.

Our results show that one may define a design score, $\Delta$, that takes on small values for sequences mounted on their true native state and large positive values for improper mounting. It is striking that the simple free energy form as in (5) can be so effective for building a reliable $\Delta$ functional. A physically appealing explanation for this is to regard the parameters in (5) as controlling both the residue composition of the designed sequence as well as indicating its expected ground state energy. The solution to a design problem will be provided by the sequence(s) that meets the composition requirements and which, when mounted on the target structure, has an energy equal to or better than the expected value. Thus, the variational approach provides a feedback mechanism for design; it is self-regulating in that no external action is required to rule out runaway solutions favoring the abundance of the most energetically favored contacts. This self-regulating mechanism also counterbalances an improper parametrization of $H$ and/or $F$, thus decreasing the sensitivity of the overall $(H-F)$ score to the detailed functional form of $\Delta$.

In conclusion, we have presented a novel procedure for tackling the direct and inverse folding problems simulta-
neously. The proposed strategy is general and ought to be applicable to the case of real proteins. We have discussed a practical implementation of the technique and have carried out rigorous testing of its efficiency in folding and design. The results are encouraging and are suggestive of the feasibility of a simple parametrization of the free energy of sequences of amino acids.

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